

Spectroscopy from the lattice: the scalar glueball and $K\pi$ scattering

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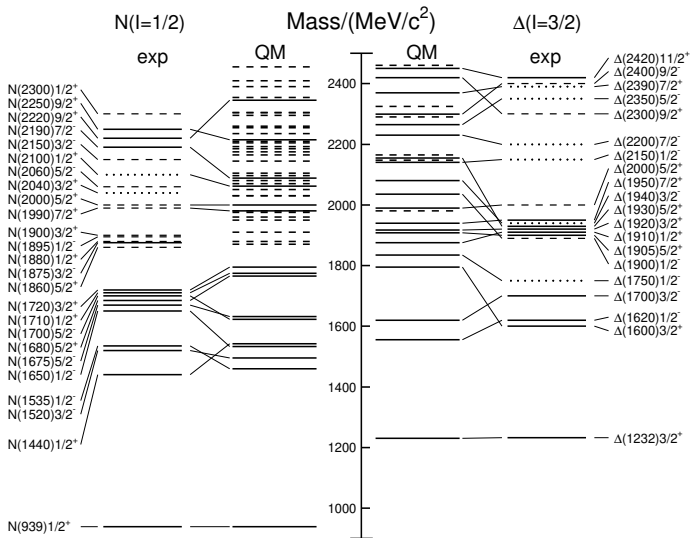
June 6, 2019

MENU 2019

Motivation

- significant experimental efforts to map out QCD spectrum: Hall B/D (JLab), COMPASS (CERN), etc.
- few states (QCD)-stable: $\pi, K, N, \Lambda, \Sigma, \Xi, \Omega$
 - most are unstable resonances
- theoretical description of states important, many poorly understood
 - eg. tetraquarks ($qq\bar{q}\bar{q}$), glueballs (bound gluons), ...
- study low-lying hadronic spectrum using lattice QCD:
 - QCD stationary states in finite-vol.
 - Lüscher: $2 \rightarrow 2$ scattering amplitudes from finite-vol. energies

Experiment vs. Quark Model

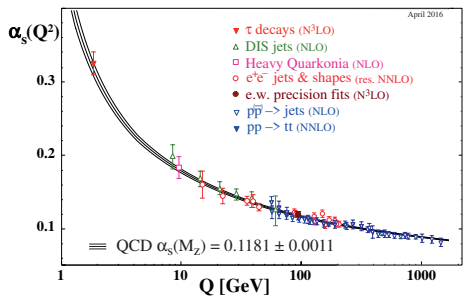


Particle Data Group

QCD on the Lattice

- strong coupling, α_s , large at low energies

→ non-perturbative



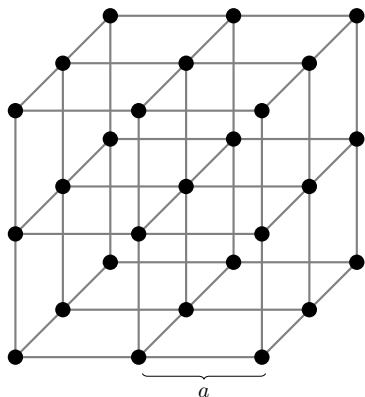
Particle Data Group

QCD on the Lattice

- strong coupling, α_s , large at low energies
 - non-perturbative
- Lattice QCD : define theory on finite, Euclidean, space-time lattice
 - lattice spacing → regulator

- observables:

$$\langle A \rangle = \frac{1}{Z} \int \mathcal{D}[\bar{\psi}, \psi, U] A e^{-S[\bar{\psi}, \psi, U]}$$



- evaluate numerically via Monte Carlo
 - ⇒ control over (many) systematic errors, etc.

Rotational Symmetry

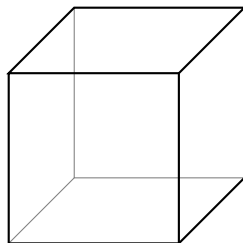
- periodic B.C. in cubic box:

⇒ J^{PC} no longer good quantum numbers

- label stationary states using irreps of cubic symmetry group O_h

- parity (u/g) and G-parity ($+/-$) where relevant

- continuum spin ID:



J	A_1	A_2	E	T_1	T_2	J	G_1	G_2	H
0	1	0	0	0	0	$\frac{1}{2}(-)^J$	1	0	0
1	0	0	0	1	0	$\frac{3}{2}(-)^J$	0	0	1
2	0	0	1	0	1	$\frac{5}{2}(-)^J$	0	1	1
3	0	1	0	1	1	$\frac{7}{2}(-)^J$	1	1	1
4	1	0	1	1	1	$\frac{9}{2}(-)^J$	1	0	2

Extracting the Finite-Volume Spectrum

- temporal. correlation matrix:

$$\begin{aligned} C_{\alpha\beta}(t) &\equiv \langle 0 | \mathcal{O}_\alpha(t + t_0) \overline{\mathcal{O}}_\beta(t_0) | 0 \rangle \\ &= \sum_n \langle 0 | \mathcal{O}_\alpha | n \rangle \langle n | \overline{\mathcal{O}}_\beta | 0 \rangle e^{-E_n t} \end{aligned}$$

- eigenvalues tend to lowest N energies

$$\begin{aligned} \lim_{t \rightarrow \infty} \lambda_n(t) &= b_n e^{-E_n t} [1 + \mathcal{O}(e^{\Delta_n t})] \\ \Rightarrow E_{\text{eff}}^n(t) &= \frac{1}{\Delta t} \ln \left(\frac{\lambda_n(t)}{\lambda_n(t + \Delta t)} \right) \end{aligned}$$

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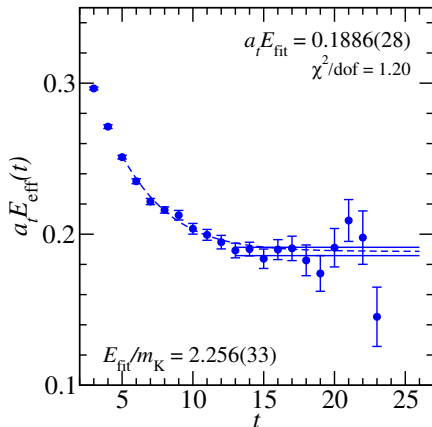
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quality of signal/excited state contamination *highly* dependant on operator basis

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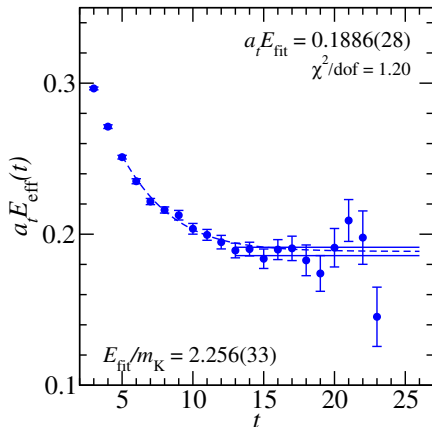
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quality of signal/excited state contamination *highly* dependant on operator basis

careful operator design is crucial

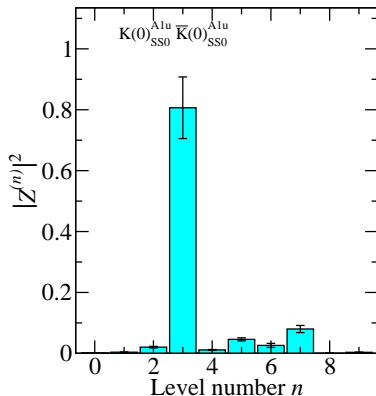
Extracting the Finite-Volume Spectrum

- 2-pt. correlation matrix:

$$\begin{aligned} C_{\alpha\beta}(t) &= \sum_n \langle 0 | \mathcal{O}_\alpha | n \rangle \langle n | \overline{\mathcal{O}}_\beta | 0 \rangle e^{-E_n t} \\ &= \sum_n Z_\alpha^{(n)} Z_\beta^{(n)*} e^{-E_n t} \end{aligned}$$

- level ID inferred from Z overlaps with *probe* operators:

$$|\Phi_j\rangle \equiv \mathcal{O}_j|0\rangle \quad \Rightarrow \quad Z_j^{(n)} = \langle \Phi_j | n \rangle$$



multiple operators can mix/overlap on any given eigenstate

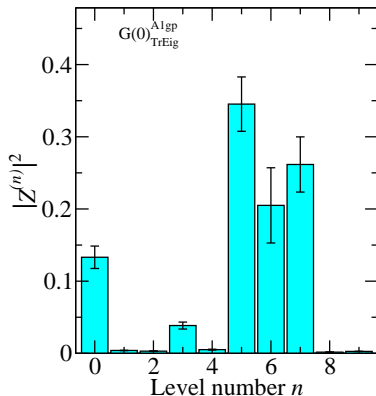
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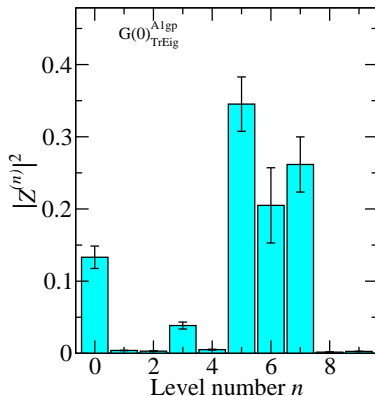
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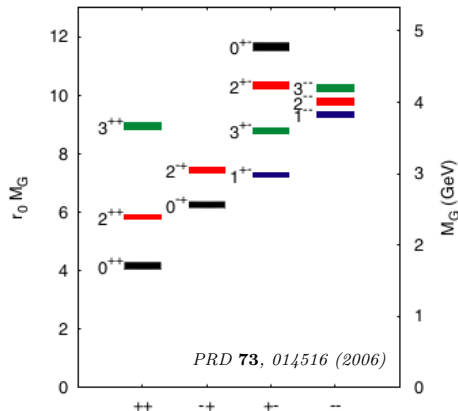


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overlaps give **qualitative** measure of mixing between states

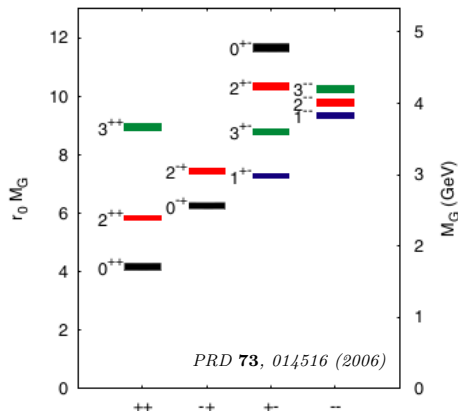
The Scalar Glueball

- glueball: hypothetical bound state of gluons
- experimental evidence elusive, light scalar candidates:
 - $f_0(1370)$, $f_0(1500)$, $f_0(1710)$
- lattice studies to date:
 - light scalar ~ 1700 MeV
 - most in pure $SU(3)$ /quenched approx. (no quark/meson mixing)



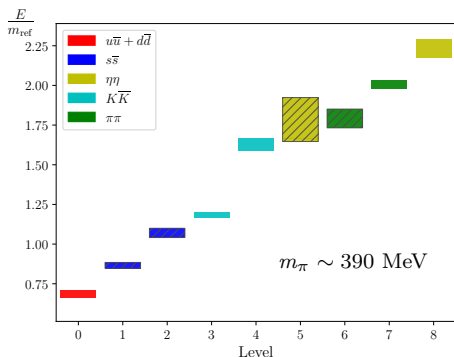
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- here: extract low-lying A_{1g}^+ spectrum with $q\bar{q}$, meson-meson, & glueball operators
 - first look (from the lattice) at mixing between glueball, $q\bar{q}$, and two-hadron states

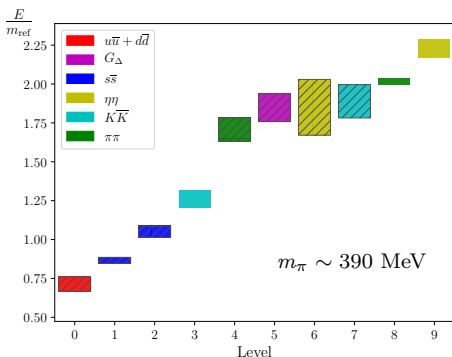


A_{1g}^+ Spectrum

$24^3 \times 128$ anisotropic lattice, $m_\pi \sim 390$ MeV, $m_K \sim 550$ MeV:



w/o glueball operator

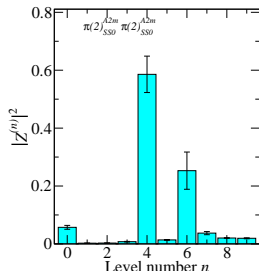
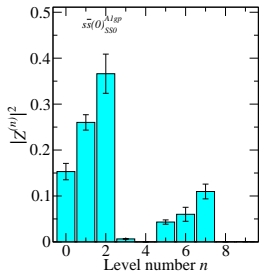
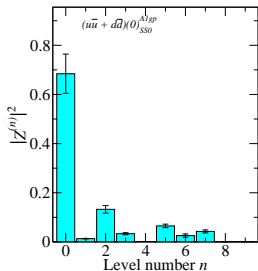
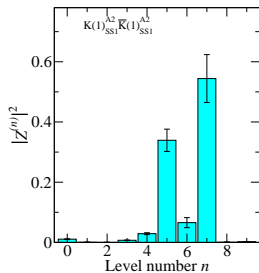
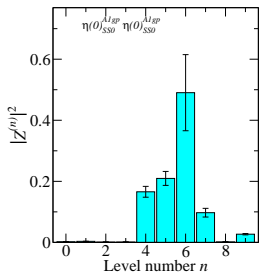
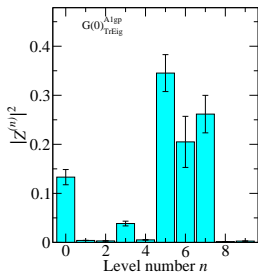


with glueball operator

hatched boxes: significant overlap with multiple operators

$$m_{\text{ref}} = 1.82m_K \sim 1 \text{ GeV}$$

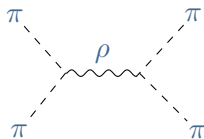
A_{1g}^+ Overlaps



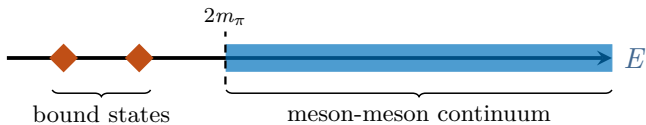
Infinite-Volume Physics

Scattering process: eg.

$$I = 1 \quad \pi\pi \rightarrow \pi\pi$$



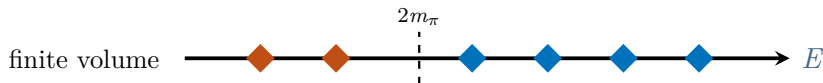
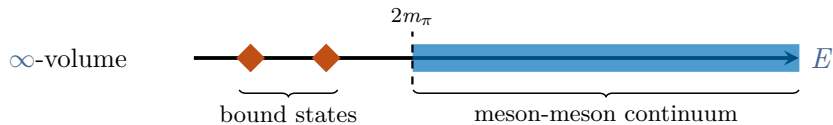
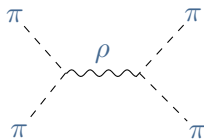
∞ -volume



Infinite-Volume Physics

Scattering process: eg.

$$I = 1 \quad \pi\pi \rightarrow \pi\pi$$



no continuum of scattering states \rightarrow how to access ∞ -vol. physics?

$$p = \frac{2\pi}{L} d$$

Lüscher Quantisation ($2 \rightarrow 2$)

Infinite volume physics from LQCD

Lüscher: Relationship between finite volume spectrum and infinite volume scattering matrix

- Quantisation condition:

$$\det[\tilde{K}^{-1} - B] = 0$$

- For each E_{cm} in spectrum, determinant gives single relation to entire scattering matrix
 - \Rightarrow Exactly solvable for single channel, single partial wave
 - \Rightarrow ℓ mixing/coupled decay channels requires parameterisation of \tilde{K} and a fit

[Nucl. Phys. B **924** 447-507 (2017)]

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↑
box matrix: known function of (E_{cm}, L)

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[Nucl. Phys. B **924** 447-507 (2017)]

Elastic $K\pi$ Scattering

- $32^3 \times 256$ anisotropic lattice,
 $m_\pi \approx 230$ MeV
- include $\ell = 0, 1, 2$ partial waves
- fit forms

$$\begin{aligned}(\tilde{K}^{-1})_{00}^{\text{ERE}} &= \frac{-1}{m_\pi a_0} + \frac{m_\pi r_0}{2} \frac{\mathbf{q}_{\text{cm}}^2}{m_\pi^2} \\(\tilde{K}^{-1})_{11}^{\text{BW}} &= \frac{6\pi E_{\text{cm}}}{g^2 m_\pi} \left(\frac{m_{K^*}^2}{m_\pi^2} - \frac{E_{\text{cm}}^2}{m_\pi^2} \right) \\(\tilde{K}^{-1})_{22}^{\text{ERE}} &= \frac{-1}{m_\pi^5 a_2}\end{aligned}$$

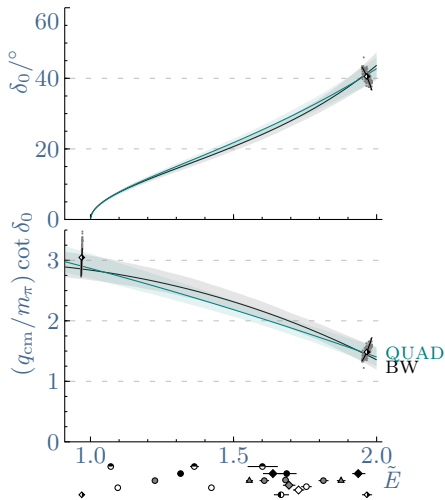
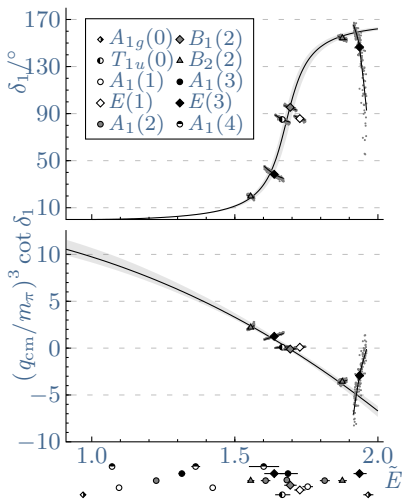
- p -wave : $K^*(892)$
- s -wave : $K_0^*(700)/\kappa$

\mathbf{d}	Λ	ℓ
$(0, 0, 0)$	A_{1g}	$0, 4, \dots$
	T_{1u}	$1, 3, \dots$
$(0, 0, n)$	A_1	$0, 1, 2, \dots$
	E	$1, 2, 3, \dots$
$(0, n, n)$	A_1	$0, 1, 2, \dots$
	B_1	$1, 2, 3, \dots$
	B_2	$1, 2, 3, \dots$
(n, n, n)	A_1	$0, 1, 2, \dots$
	E	$1, 2, 3, \dots$

[Nucl.Phys. B **932** 29-51 (2018)]

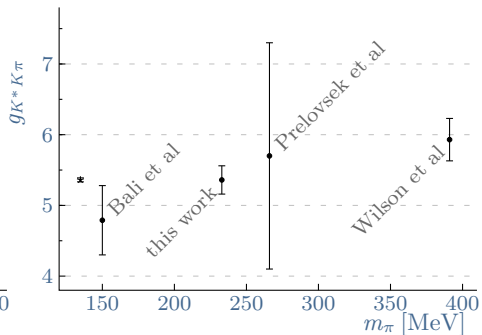
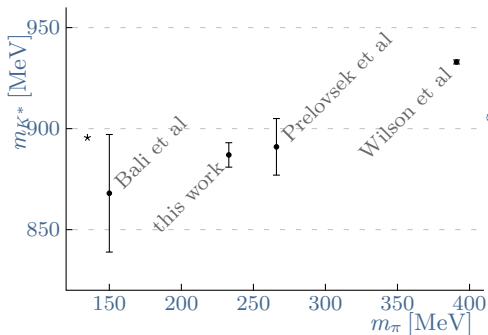
$K\pi$ Scattering Phase-Shifts

- plots of p -wave and s -wave phase shift ($\tilde{E} = (E_{\text{cm}} - m_K)/m_\pi$)
- d -wave contributes negligibly



Decay of $K^*(892)$

- comparison with (some) lattice calculations of $K^*(892)$ resonance parameters
- phenomenological values shown with asterisks



Conclusions/Future work - Scattering

- goal: calculate scattering phase-shifts/resonance parameters from LQCD
- meson-meson scattering at a mature stage
 - see B. Hörz's talk today
 - physical point calculations feasible - large volumes required
- meson-baryon scattering now possible: eg. $N\pi$: PRD **97**, 014506 (2018)
 - see J. Bulava's plenary this morning
 - baryon-baryon at light quark masses on the horizon
- extension of formalism to include three-body channels underway
 - see M. Mai's talk on Mon.

Conclusions/Future work - Spectroscopy

- goal: *qualitative* description of resonance spectrum
- high computational cost for large operator bases/volumes
 - careful operator design crucial
- $\bar{q}q$ states straightforward but many interesting states not well described
 - ⇒ hybrids
 - ⇒ molecular states
 - ⇒ ...
- influence of tetraquark operators: κ and $a_0(980)$
 - see D. Darvish's talk on Tues.

Meson and Baryon elemental operators



- q 's = smeared, displaced quark fields

$$\bar{\Phi}_{\alpha\beta;ij}^{AB}(\mathbf{p}, t) = \sum_{\mathbf{x}} e^{i\mathbf{p}\cdot(\mathbf{x} + \frac{1}{2}(\mathbf{d}_\alpha + \mathbf{d}_\beta))} \delta_{ab} \bar{q}_{b\beta j}^B(\mathbf{x}, t) q_{a\alpha i}^A(\mathbf{x}, t)$$

$$\bar{\Phi}_{\alpha\beta\gamma;ijk}^{ABC}(\mathbf{p}, t) = \sum_{\mathbf{x}} e^{i\mathbf{p}\cdot\mathbf{x}} \varepsilon_{abc} \bar{q}_{c\gamma k}^C(\mathbf{x}, t) \bar{q}_{b\beta j}^B(\mathbf{x}, t) \bar{q}_{a\alpha i}^A(\mathbf{x}, t)$$

- group-theory projections onto irreps of lattice symmetry group

$$\bar{M}_l(t) = c_{\alpha\beta}^{(l)*} \bar{\Phi}_{\alpha\beta}^{AB}(t) \quad \bar{B}_l(t) = c_{\alpha\beta\gamma}^{(l)*} \bar{\Phi}_{\alpha\beta\gamma}^{ABC}(t)$$

More exotic operators

- tetraquarks: $3 \times 3 \times \bar{3} \times \bar{3} = 1 + 1 + 8 + 8 + \dots$

$$\begin{aligned} \bar{\Phi}_{\alpha\beta\mu\nu;ijkl}^{ABCD\pm}(\mathbf{p}, t) &= \sum_{\mathbf{x}} e^{i\mathbf{p}\cdot\mathbf{x}} (\delta_{ab}\delta_{cd} \pm \delta_{ad}\delta_{bc}) \\ &\quad \times \bar{q}_{d\nu l}^D(\mathbf{x}, t) q_{c\mu k}^C(\mathbf{x}, t) \bar{q}_{b\beta j}^B(\mathbf{x}, t) q_{a\alpha i}^A(\mathbf{x}, t) \end{aligned}$$

→ preprint coming soon: κ & $a_0(980)$ resonances

More exotic operators

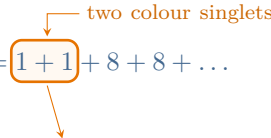
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two colour singlets!

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- scalar glueball: (purely gluonic - no quark fields)

$$G_{\Delta} \sim \text{Tr } \tilde{\Delta}, \quad \tilde{\Delta} \equiv \text{covariant laplacian}$$

→ see PRD **88**, 014511 (2013)

S -wave $K\pi$ amplitude: $K_0^*(700)/\kappa$

$$\frac{m_{K^*}}{m_\pi} = 3.808(18), \quad g = 5.33(20), \quad m_\pi a_0 = -0.353(25),$$
$$m_\pi^5 a_2 = -0.0013(68), \quad \chi^2/\text{dof} = 1.42$$

- based on LO ERE, $m_\pi a_0 < 0$ suggests virtual bound state
- however, NLO parameters give $1 - 2r_0/a_0 = -8.9(2.4)$ which must be > 0 for a (real or virtual) bound state
- zeros of $\mathbf{q}_{\text{cm}} \cot \delta_0 - i\mathbf{q}_{\text{cm}}$: $m_R/m_\pi = 4.66(13) - 0.87(18)i$
 - consistent with BW fit
- better energy resolution & careful analytic continuation required

[Nucl.Phys. B932 (2018)]