

The $\pi N \sigma$ term from pionic atoms

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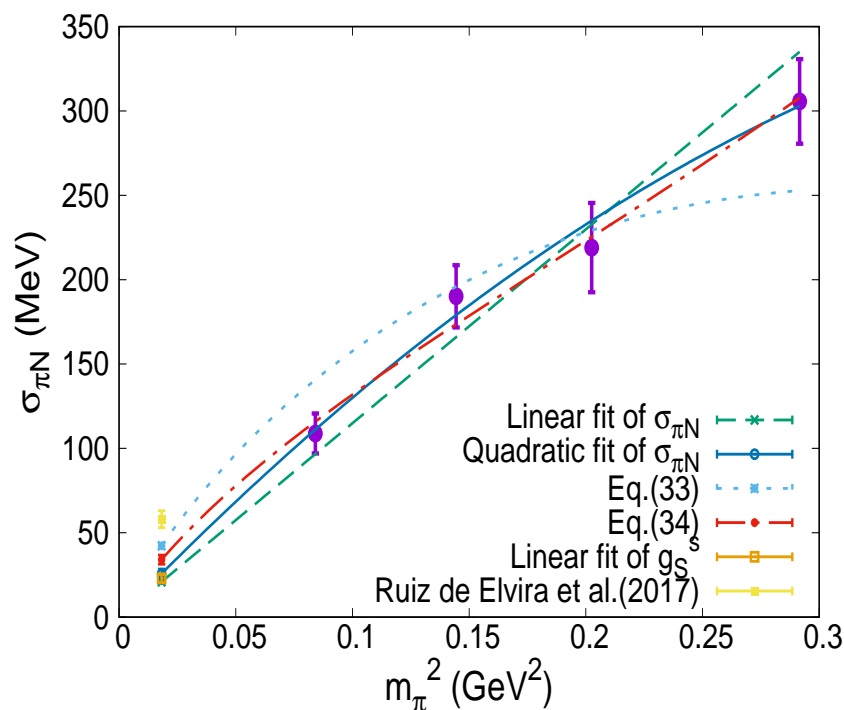
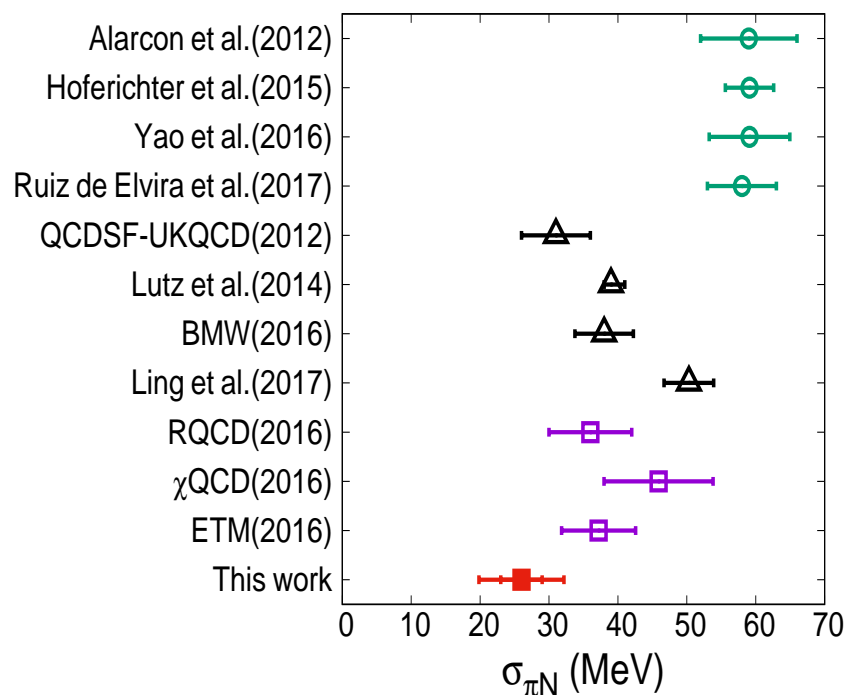
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- Partial restoration of chiral symmetry from pionic atoms. Last update: **NPA 928 (2014) 128.**
- Extracting $\sigma_{\pi N}$ from pionic atom data: **PLB 792 (2019) 340.**
- Comparison with other methods
 - (i) πN dynamics near threshold ($b_0 \approx 0$)
(Hoferichter's plenary talk): $\sigma_{\pi N} \sim 60$ MeV
 - (ii) LQCD calculations: $\sigma_{\pi N} \sim 30$ MeV.

The pion-nucleon σ term

$$\sigma_{\pi N} = \frac{\bar{m}_q}{2m_N} \Sigma_{u,d} \langle N | \bar{q}q | N \rangle, \quad \bar{m}_q = \frac{1}{2}(m_u + m_d)$$

records the contribution of explicit chiral symmetry breaking to the nucleon mass m_N arising from the non-zero value of the u and d quark masses in QCD.



various calcs. of $\sigma_{\pi N}$

chiral extraps. of $\sigma_{\pi N}$

N. Yamanaka et al. (JLQCD) PRD 98 (2018) 054516

Partial restoration of chiral symmetry in/from pionic atoms

Update: Friedman-Gal, NPA 928 (2014) 128

Hadronic (h) atom scenarios

h	Re V_{opt}	Im V_{opt}	comments
π^-	vol. repulsion surf. attraction	moderate	excellent data well underst.
Σ^-	vol. repulsion tail attraction	moderate	limited data poorly underst.
K^-	attractive how deep ?	absorptive 1N vs. 2N	good data open problems
\bar{p}	??	absorptive	excellent data

Optical model analyses of hadronic atom data

- Handle large data sets across periodic table.
- Identify characteristic entities, thereby linking microscopic approaches to experiments.

Tools of the trade: optical potential variants

- Make V_{opt} functional of the nuclear density ρ .
- Respect the low-density limit $V_{\text{opt}}(\rho) \rightarrow t_{hN} * \rho$.
- For pions, consider $\rho_n - \rho_p$ dependence of b_1 using $r_n - r_p \approx \gamma \frac{N-Z}{A} + \delta$ with $\gamma \approx 1.0 \pm 0.1$ fm.
- Introduce self consistently medium effects, particularly subthreshold hN kinematics.

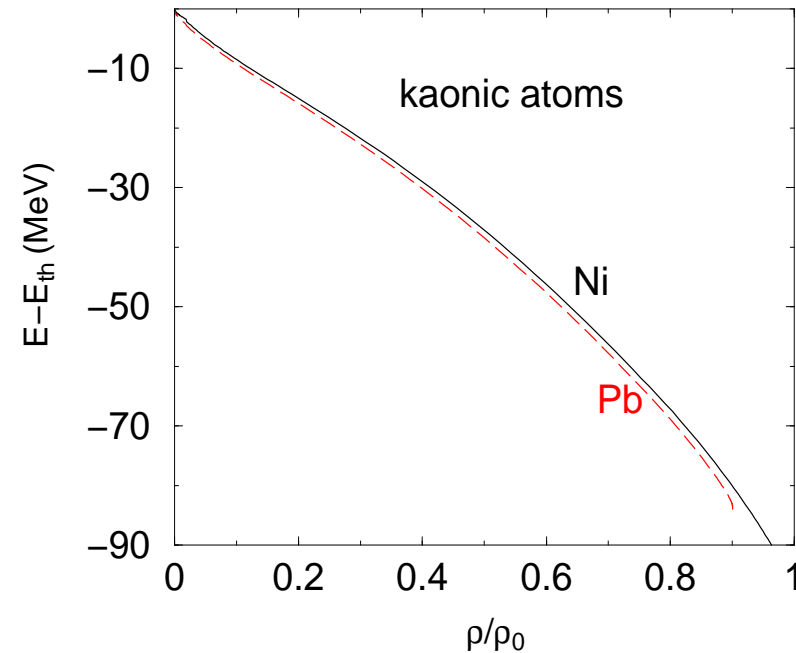
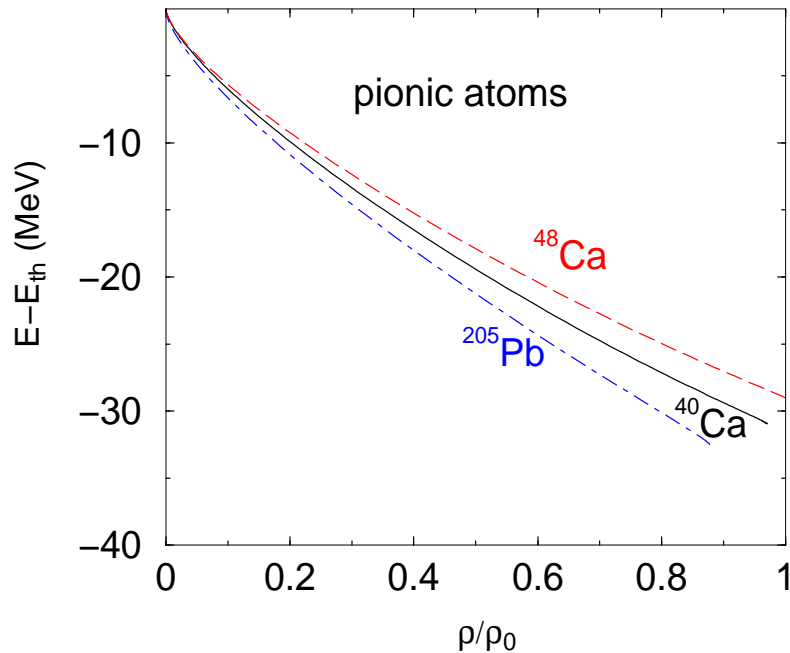
Self-consistency in mesic-atom & nuclear calculations

Cieplý-Friedman-Gal-Gazda-Mareš, PLB 702 (2011) 402

$$s_{hN} = (\sqrt{s_{\text{th}}} - B_h - B_N)^2 - (\vec{p}_h + \vec{p}_N)^2 < s_{\text{th}}$$

$$\sqrt{s_{hN}} \rightarrow E_{\text{th}} - B_N - B_h - \xi_N \frac{p_N^2}{2m_N} - \xi_h \frac{p_h^2}{2m_h}$$

$$\xi_{N(h)} = \frac{m_{N(h)}}{(m_N + m_h)} \quad \frac{p_h^2}{2m_h} \sim -V_h - B_h$$



in-medium π^- energy shift

in-medium K^- energy shift

In-medium hN amplitudes

Friedman-Gal-Mareš, PLB 725 (2013) 334

Cieplý-Friedman-Gal-Mareš, NPA 925 (2014) 126

- KG equation and self-energies:

$$[\nabla^2 + \tilde{\omega}_h^2 - m_h^2 - \Pi_h(\omega_h, \rho)] \psi = 0$$

$$\tilde{\omega}_h = \omega_h - i\Gamma_h/2, \quad \omega_h = m_h - B_h$$

$$\Pi_h(\omega_h, \rho) \equiv 2\omega_h V_h = -4\pi \frac{\sqrt{s}}{m_N} f_{hN}(\sqrt{s}, \rho) \rho$$

- Pauli blocking (Waas-Rho-Weise NPA 617 (1997) 449):

$$f_{hN}^{\text{WRW}}(\sqrt{s}, \rho) = \frac{f_{hN}(\sqrt{s})}{1 + \xi(\rho)(\sqrt{s}/m_N)f_{hN}(\sqrt{s})\rho}, \quad \xi(\rho) = \frac{9\pi}{4p_F^2} I(\tilde{\omega}_h)$$

- \sqrt{s} : $\Lambda^*(1405) \Rightarrow f_{K-N}(\sqrt{s})$, $N^*(1535) \Rightarrow f_{\eta N}(\sqrt{s})$.

In medium \Rightarrow go subthreshold: $\delta\sqrt{s} = \sqrt{s} - \sqrt{s_{\text{th}}}$

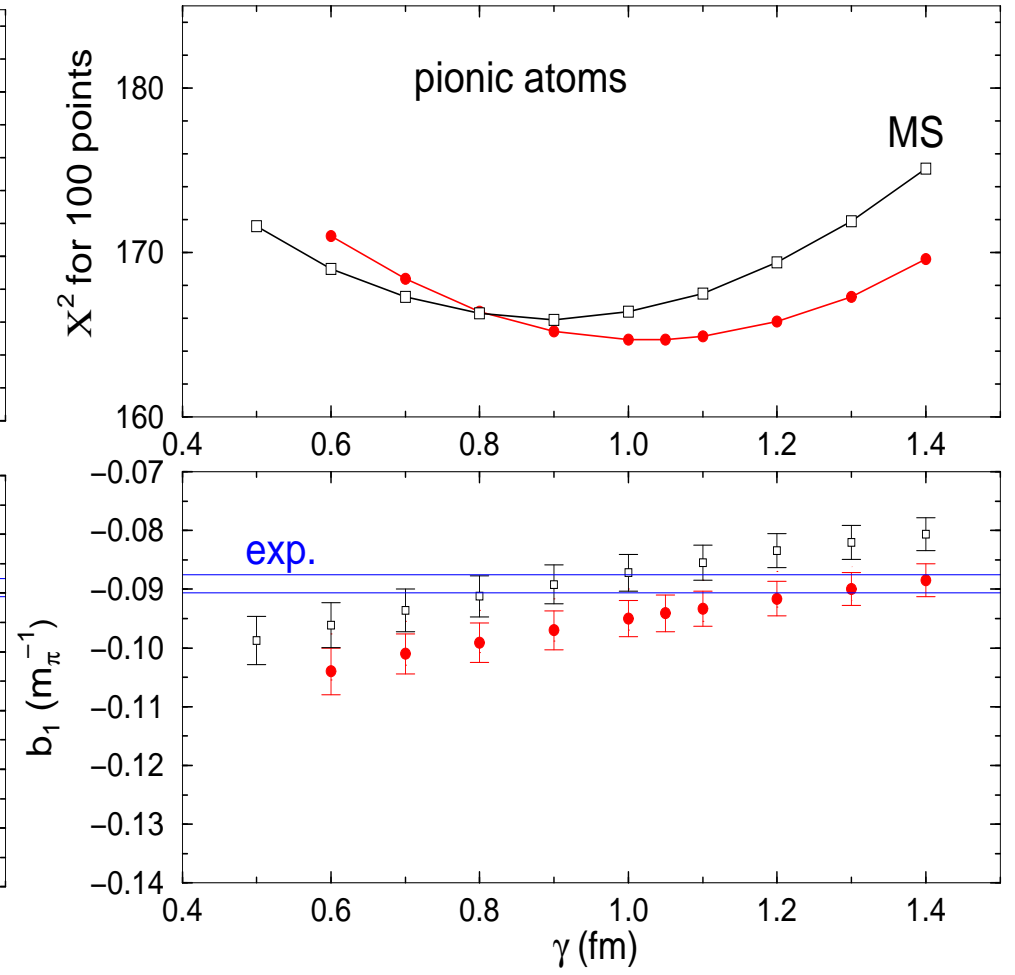
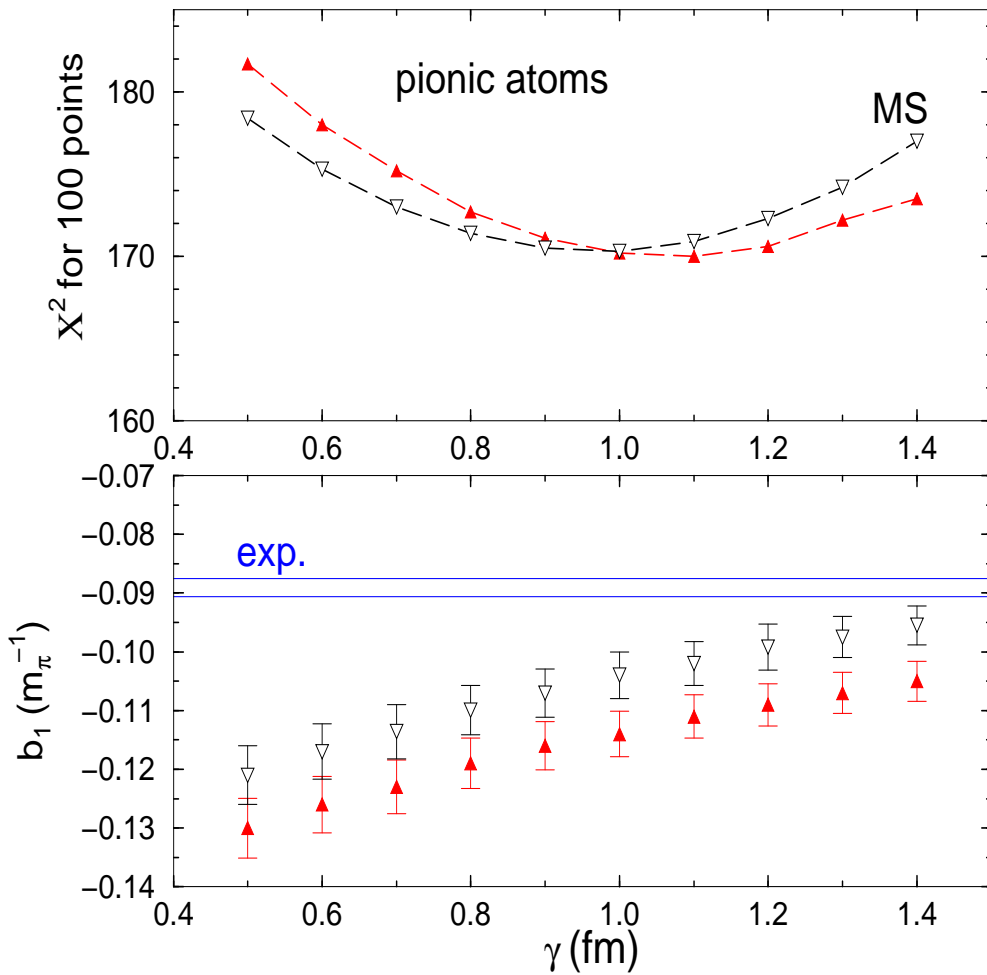
$$\delta\sqrt{s} \approx -B_N \frac{\rho}{\rho_0} - \xi_N B_h \frac{\rho}{\rho_0} - \xi_N T_N \left(\frac{\rho}{\rho_0}\right)^{2/3} + \xi_h \text{Re} V_h(\sqrt{s}, \rho)$$

- **A self-consistency cycle in $\delta\sqrt{s}$ for given ρ .**

Pion-nucleus optical potential

Ericson-Ericson form (1966)

- $2\mu V_{\text{opt}}(r) = q(r) + \vec{\nabla} \cdot \left(\frac{\alpha_1(r)}{1 + \frac{1}{3}\xi\alpha_1(r)} + \alpha_2(r) \right) \vec{\nabla}$
with s-wave $q(r)$ and p-wave $\alpha(r)$.
- $q(r) \sim b_0[\rho_n(r) + \rho_p(r)] + b_1[\rho_n(r) - \rho_p(r)] + 4B_0\rho_n(r)\rho_p(r)$, $b_0 \rightarrow b_0 - \frac{3}{2\pi}(b_0^2 + 2b_1^2)p_F$
- On-shell values from π^- -H & π^- -D atoms (PSI):
 $b_0^{\text{free}} = 0.0076(31) m_\pi^{-1}$ $b_1^{\text{free}} = -0.0861(9) m_\pi^{-1}$
LO χ limit: $b_0^{\text{TW}} = 0$, $b_1^{\text{TW}} = -\frac{\mu_{\pi N}}{8\pi f_\pi^2} = -0.079 m_\pi^{-1}$
- GMOR: $\frac{f_\pi^2(\rho)}{f_\pi^2} = \frac{\langle \bar{q}q \rangle_\rho}{\langle \bar{q}q \rangle_0} \simeq 1 - \frac{\sigma\rho}{m_\pi^2 f_\pi^2}$
 $\Rightarrow \sim b_1^{-1}(\rho)$
- $\alpha_1(\mathbf{r})$: $b_0 \rightarrow c_0$, $b_1 \rightarrow c_1$ $\alpha_2(\mathbf{r})$: $B_0 \rightarrow C_0$

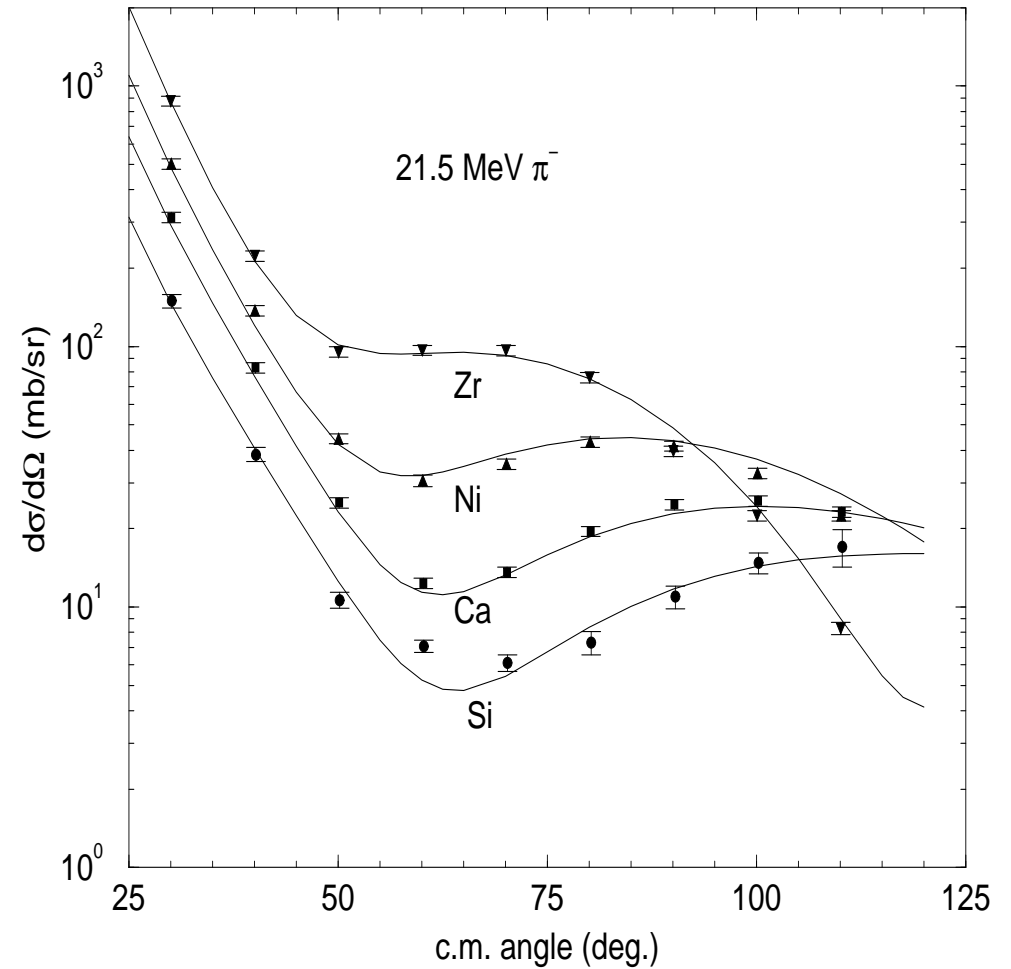
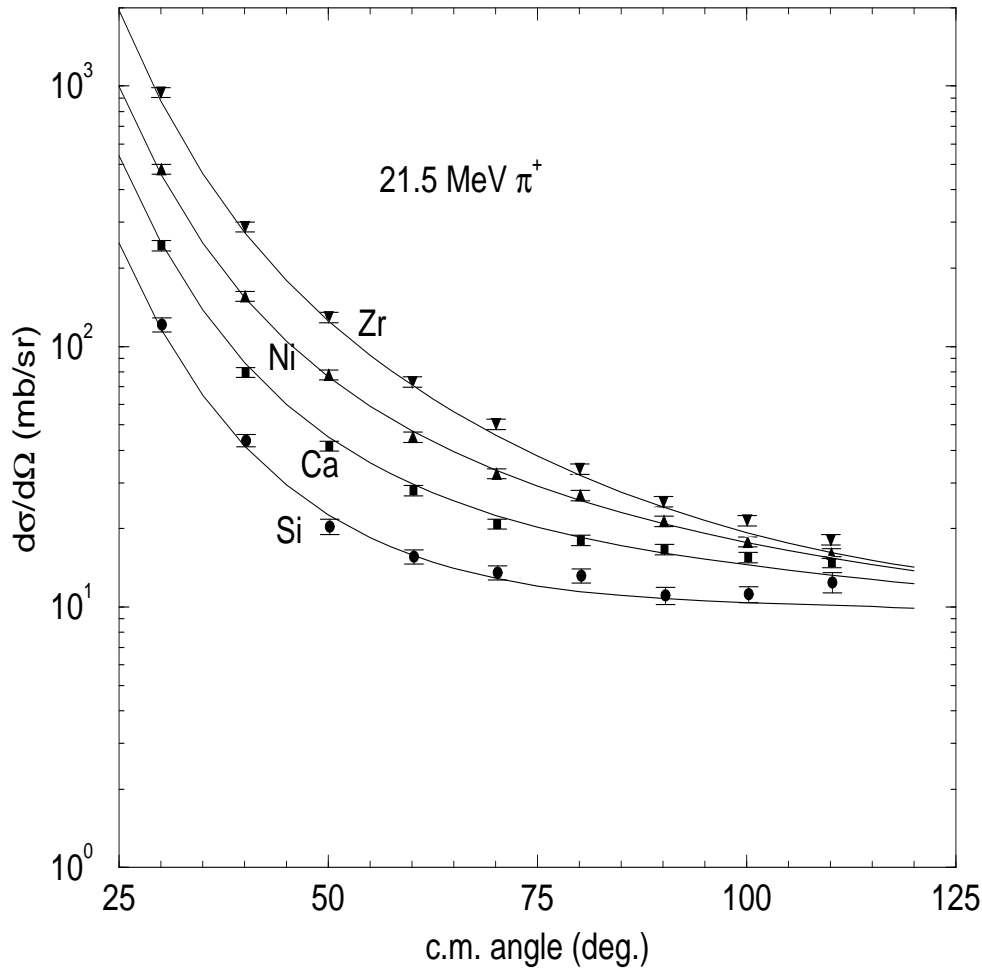


empirical $b_0(E)$ using
 b_1 independent of ρ

empirical $b_0(E)$ using
 $b_1(\rho)$ with $\sigma_{\pi N} = 50$ MeV

γ dependence of fits to 100 data points, Ne to U
 More precise determination of b_1 than thru DBS

E. Friedman, A. Gal, NPA 928 (2014) 128



E. Friedman et al., PRL 93 (2004) 122302, PRC 72 (2005) 034609
PSI results reproduced with $b_1(\rho)$ ansatz (Weise, 2000)

$$b_1(\rho) = -\frac{\mu_{\pi N}}{8\pi f_\pi^2(\rho)}, \quad \frac{f_\pi^2(\rho)}{f_\pi^2} = \frac{\langle \bar{q}q \rangle_\rho}{\langle \bar{q}q \rangle_0} \simeq 1 - \frac{\sigma\rho}{m_\pi^2 f_\pi^2}, \text{ applied for } \sigma=50 \text{ MeV}$$

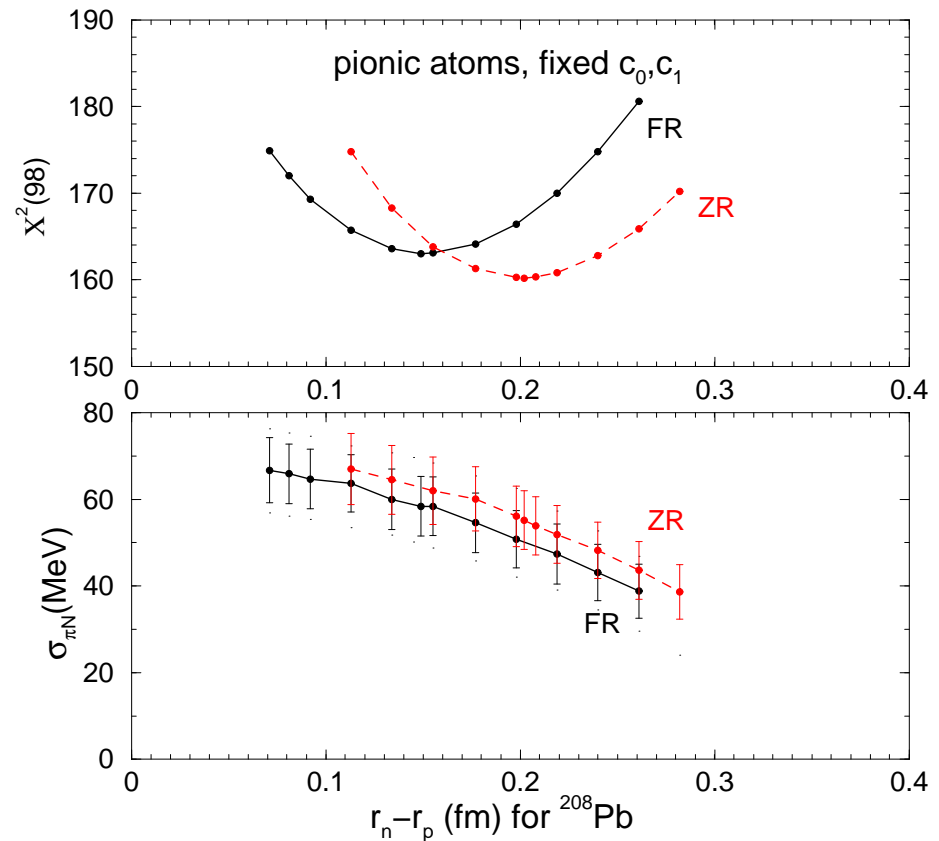
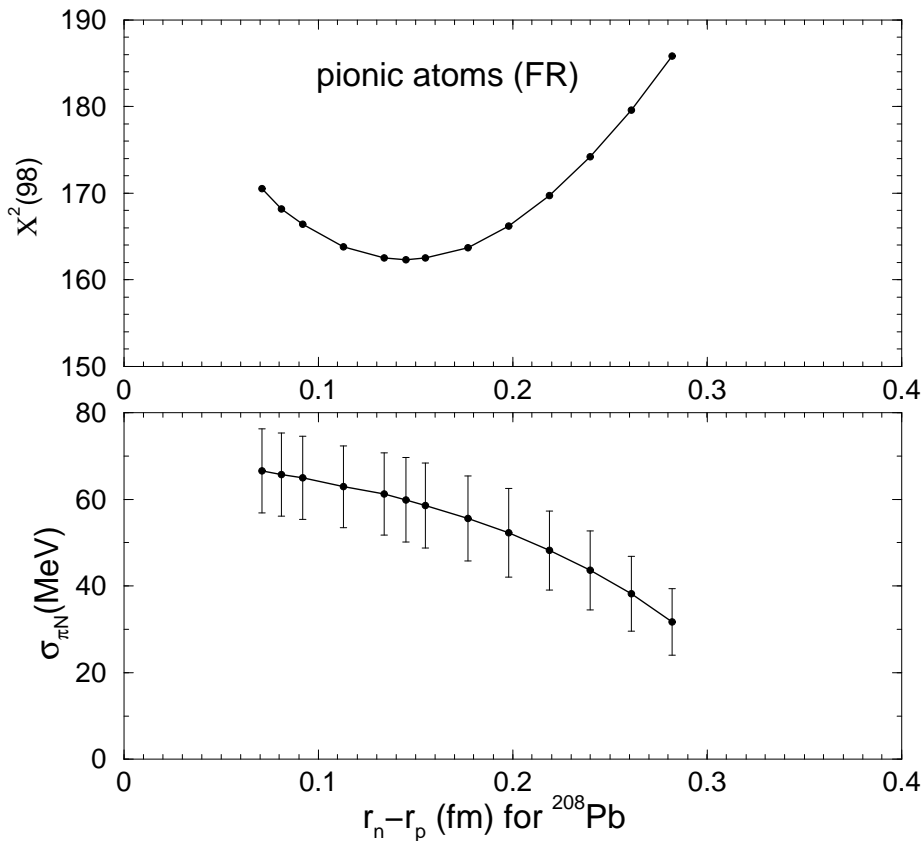
Consistency between π^- atoms & π^\pm scattering deductions.

Constraining the pion-nucleon σ term from pionic atoms

Recent: Friedman-Gal, PLB 792 (2019) 340

Fitting $\sigma_{\pi N}$ to π^- atom data

≈ 100 data point, shifts and widths, across the periodic table,
 ≈ 10 are from deeply bound pionic atoms.



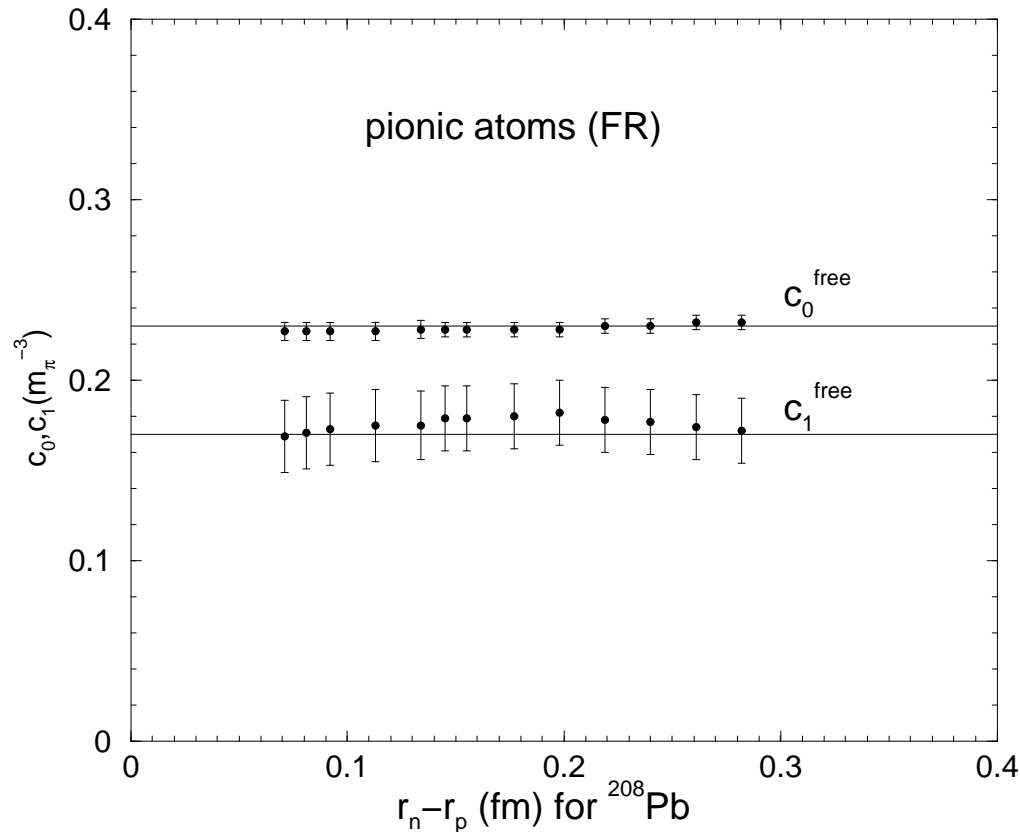
Up: χ^2 , Low: $\sigma_{\pi N}$
6 parameters

Finite & Zero range p waves
4 parameters

best-fit $\sigma_{\pi N} = 57 \pm 7 \text{ MeV}$

Discussion & Summary

Stability to p-wave potential parameters



Horizontal lines mark the SAID free-space values of the πN scattering volumes.

Resulting $\sigma_{\pi N}$ is robust to fit details.

- Corrections beyond leading density are a few percent.
- Our $\sigma_{\pi N} = 57 \pm 7 \text{ MeV}$ agrees with Hoferichter et al. PRL 115 (2015) 092301 & PLB 760 (2016) 74 value $\sigma_{\pi N} = 59.1 \pm 3.5 \text{ MeV}$ which depends primarily on extrapolating b_0 to the subthreshold Cheng-Dashen point.
- Note that the model dependence of b_0 is fairly large compared to that of b_1 upon which our pionic-atom determination relies.
- Need to improve LQCD derivations...

Thanks for your attention!